

Introduction to topological order and topological quantum computation

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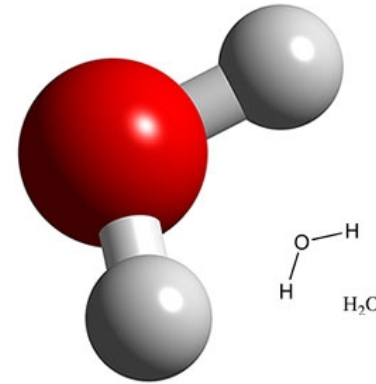
Outline

1. Introduction: phase transitions and order.
2. The Landau symmetry breaking description.
3. The topological order.
4. Signs of topological order: anyons.
5. Topological quantum computation.

Introduction: Phase transitions and the concept of order

The principle of emergence

Matter is formed by atoms/molecules.



Different organizations



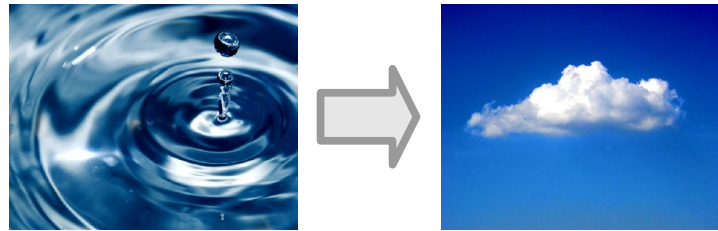
Different collective properties



Each of these organizations is called **phase**.

Phase Transitions

A **phase transition** is the transformation of a thermodynamic system from one **phase** to another.



In a **phase transition**, the **free energy** or some of its derivatives diverge.

First order PTs
involve a latent heat



Second order PTs
do not involve a latent heat



Quantum phase transitions

A **quantum phase transition** is a phase transition between two **quantum phases** at zero temperature. They can only be accessed by varying a physical parameter.

We will identify any point of **nonanalyticity** in the **ground state** energy of the infinite system as a **quantum phase transition**.

Example: The **quantum Ising model**.

$$H_{\text{Ising}} = -\frac{1}{2} \sum_l (\sigma_l^x \sigma_{l+1}^x + \lambda \sigma_l^z)$$

Quantum paramagnet phase

$$|GS\rangle = \prod_i |\uparrow\rangle_i$$

$$\langle GS | \sigma_i^x \sigma_j^x | GS \rangle \sim e^{-|x_i - x_j|/\xi}$$

Quantum ferromagnet phase

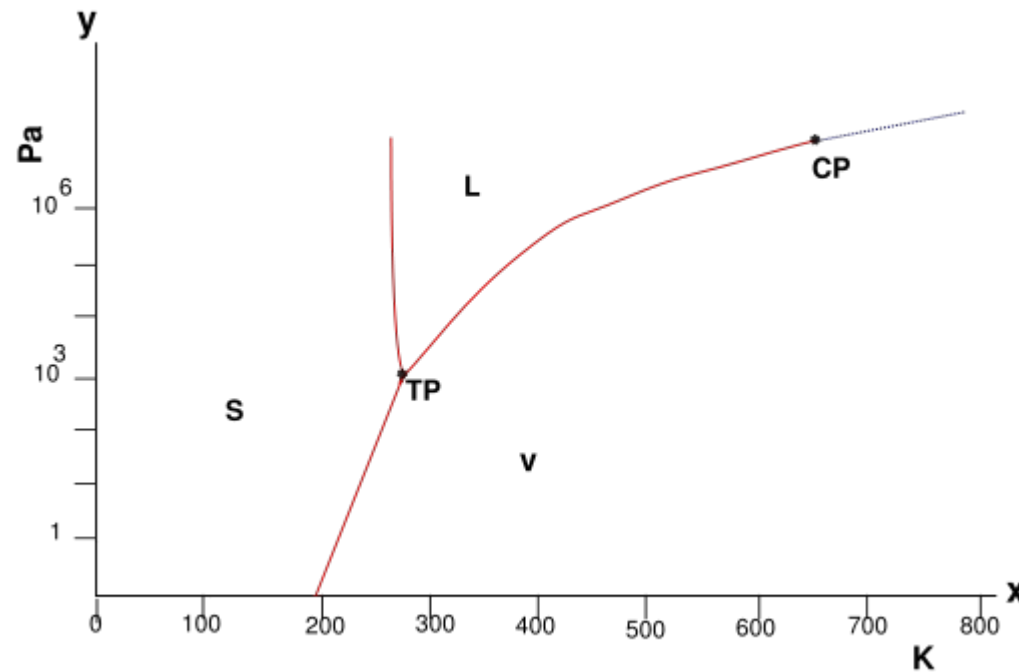
$$|GS\rangle = \prod_i |\leftarrow\rangle_i$$

$$\lim_{|x_i - x_j| \rightarrow \infty} \langle GS | \sigma_i^x \sigma_j^x | GS \rangle \sim N_0^2$$

The concept of order

The concept of **order** involves **phase transition**.

Two many particle **states** are in the same **order** if we can **smoothly** change one state into the other without encountering a **phase transition**.



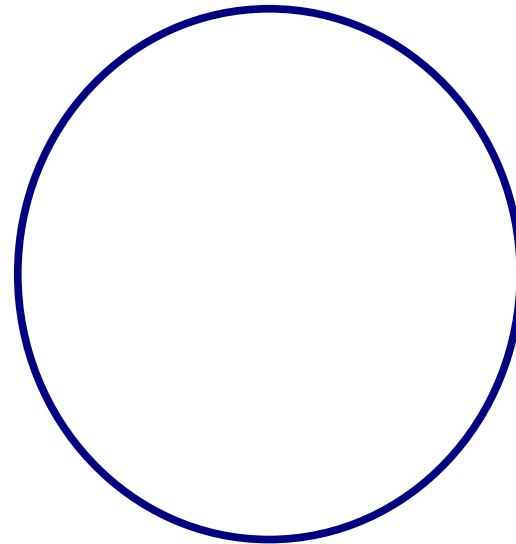
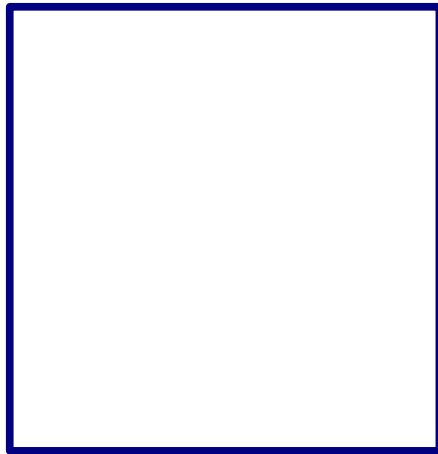
How many **orders** we have in the water **phase** diagram?

What makes two orders really different?

Landau symmetry breaking theory

Order and symmetry

Order and *symmetry* are complementary concepts.



“More order means less symmetry”

Order and symmetry

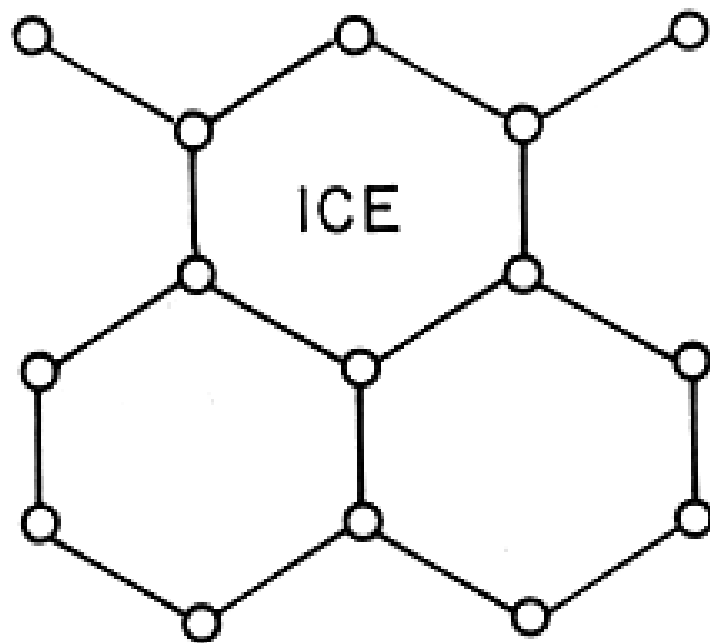
Order and *symmetry* are complementary concepts.



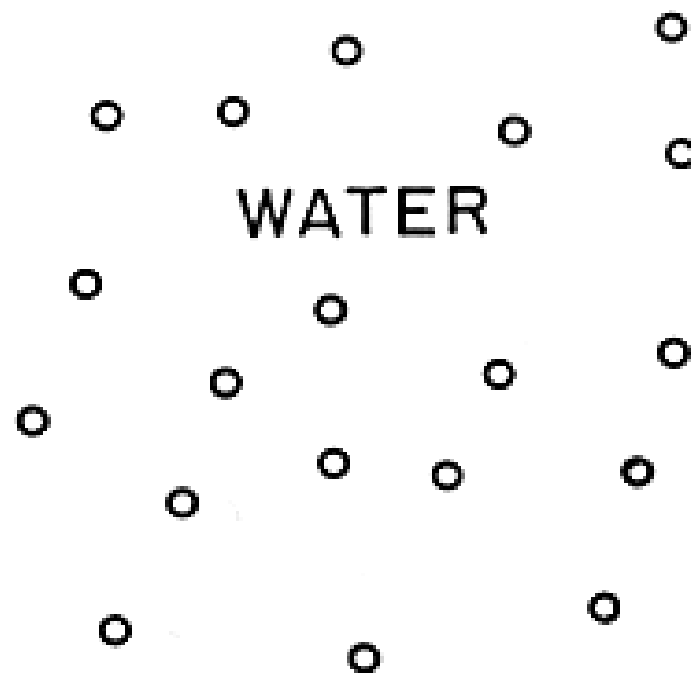
“More order means less symmetry”

Order and symmetry

Which phase is more symmetric?



A

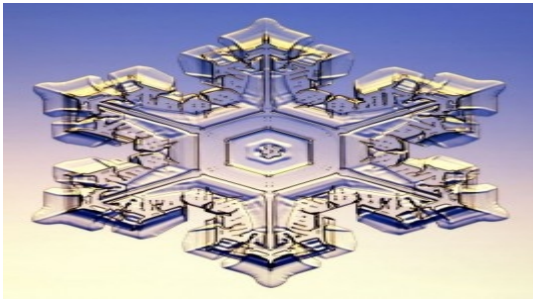


B

Landau symmetry breaking theory

The Landau symmetry breaking theory relates **order** with **symmetry**.

Different orders correspond to different symmetries



Discrete translational invariance



Continuous translational invariance



Symmetry breaking

Phonons are the Goldstone bosons associated to the symmetry breaking.

Order parameter

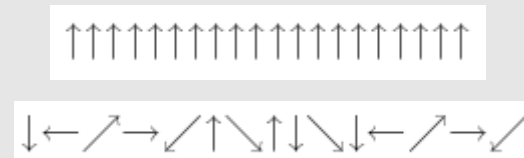
The **order parameter** is a quantity that:

- takes one value in one **phase** and another one in the other to **distinguish** in which phase the system is.
- **measures** of the **degree of order** in a system; the extreme values are 0 for total **disorder** and 1 for complete **order**.

Example

In the **Heisenberg model**, the **magnetization** is 0 in the **paramagnetic** phase and 1 in the **ferromagnetic** one.

$$H = - \sum_i \vec{S}_i \vec{S}_{i+1}$$



Which is the broken **symmetry**?

Landau's paradigm

The Landau's paradigm ingredients:

- the relation between **symmetry** (group theory) and **order**.
- the origin of many **gapless excitations**: phonons, spin waves,...
- the **order parameter**.

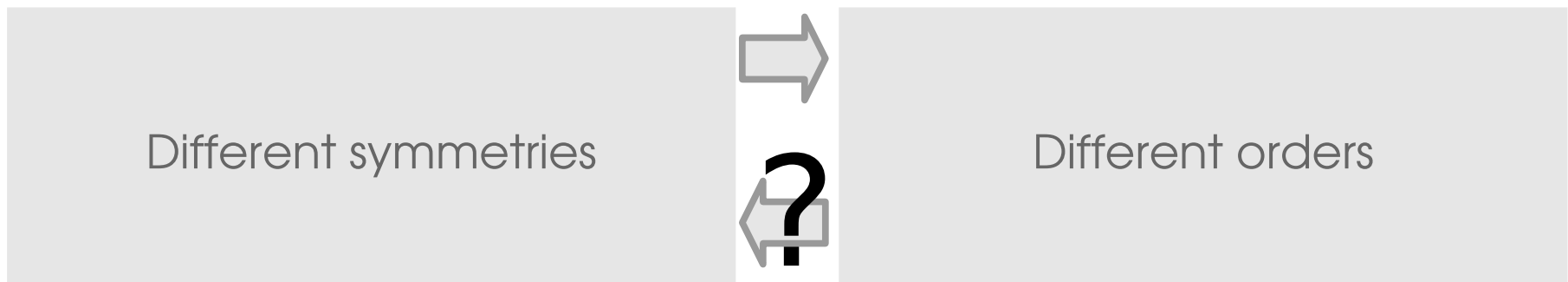
The **Landau theory** has been very successful:

- we can **classify** all of the 230 different kinds of crystals that exists in 3 dimensions.
- by determining how **symmetry** changes across a 2nd order **phase transition**, we can obtain its **critical** properties.

Beyond Landau's paradigm: topological order

Beyond Landau's paradigm

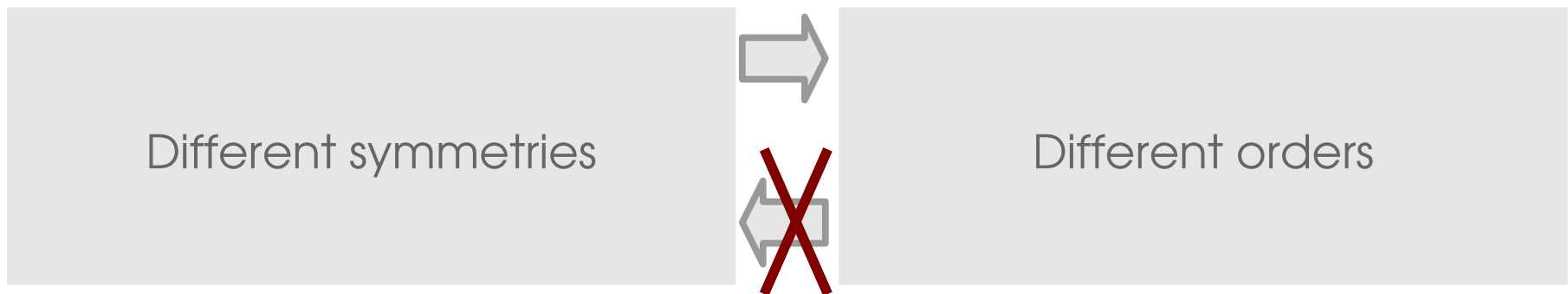
Can two **different orders** have the same **symmetry**?



Beyond Landau's paradigm

Can two **different orders** have the same **symmetry**?

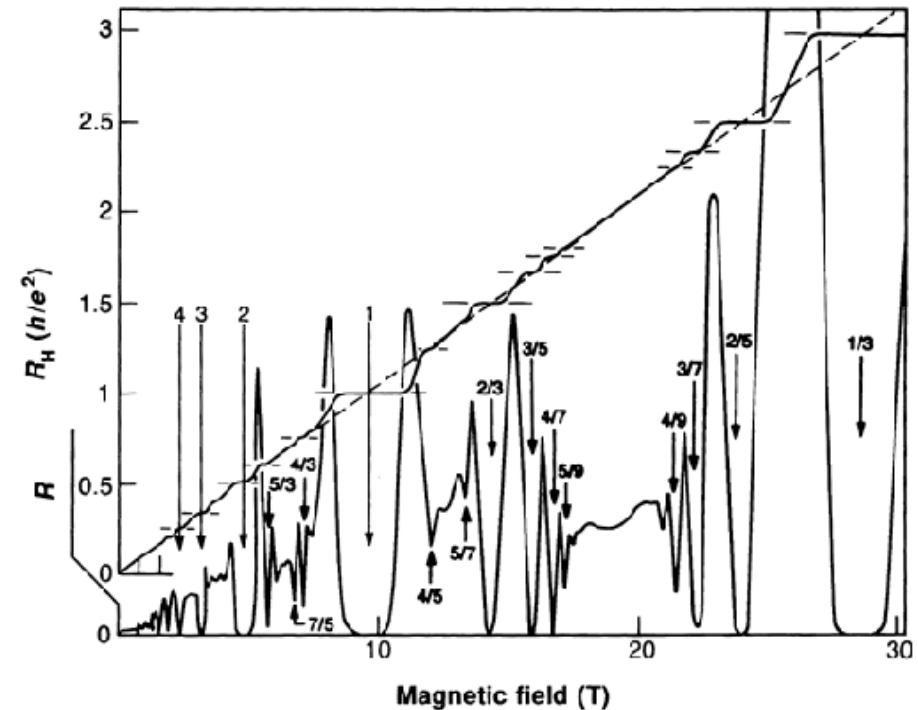
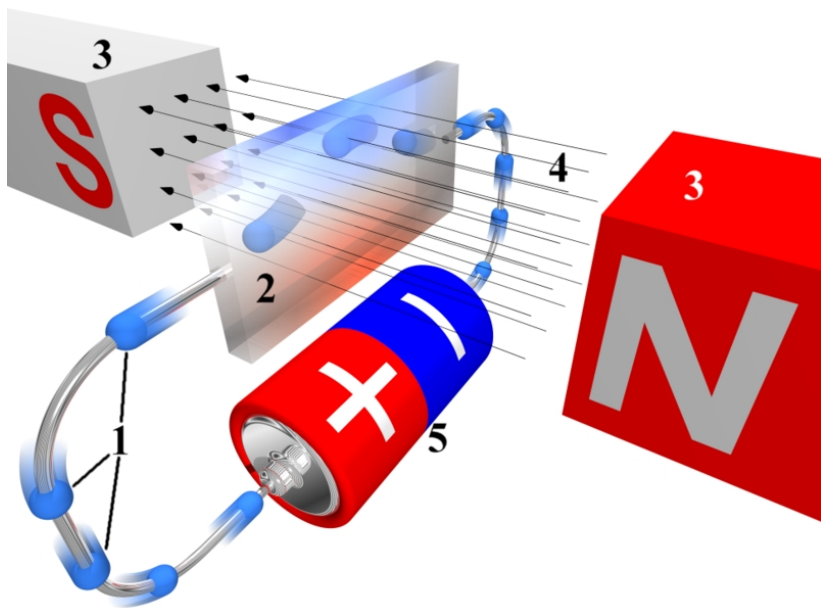
yes!



The Fractional Quantum Hall Effect (FQHE)!

The Fractional Quantum Hall Effect

In 1980, Tsui *et al.* measured the **transversed resistivity** of a **2D electron gas** (at interfaces between semiconductors) under a strong **transverse magnetic field** at very **low temperatures**.



They observed the formation of **plateaus**, where the **Hall conductivity** took values $1/3, 5/2, 2/3$ of e^2/h .

The Laughlin wave function

The [Laughlin wave function](#) was postulated by Robert B. Laughlin in 1983 as the [ground state](#) of the [FQHE](#). It is defined by

$$\Psi_L^{(m)}(z_1, \dots, z_n) \sim \prod_{i < j} (z_i - z_j)^m e^{-\sum_{i=1}^N |z_i|^2 / 2}$$

where $z_i = x_i + i y_i$.

It takes into account:

- the [single-electron](#) Hilbert space is restricted to the subspace of zero energy (lowest Landau level).
- The [statistics](#) of the particles.
- The [repulsive interaction between](#) electrons.

The Laughlin wave function

Why the Laughlin state is so special?

$$\Psi_L^{(m)}(z_1, \dots, z_n) \sim \prod_{i < j} (z_i - z_j)^m e^{-\sum_{i=1}^N |z_i|^2 / 2}$$

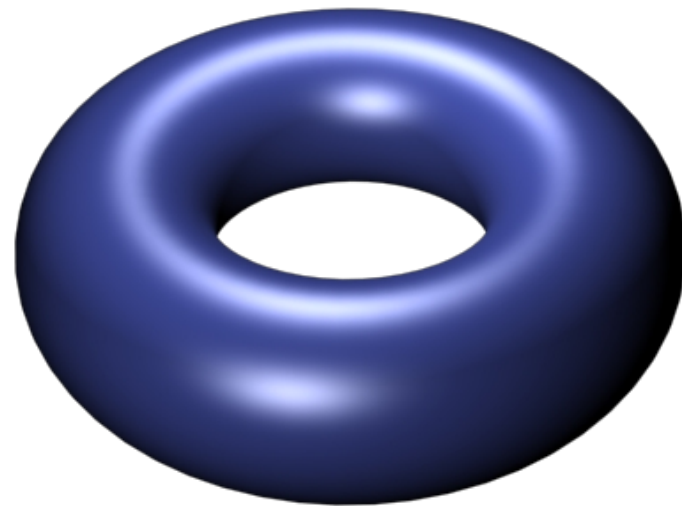
This state is completely disorderd and symmetry is not broken.

States with a different m share the same symmetry and we cannot change one into another without finding a quantum phase transition.

The topological order

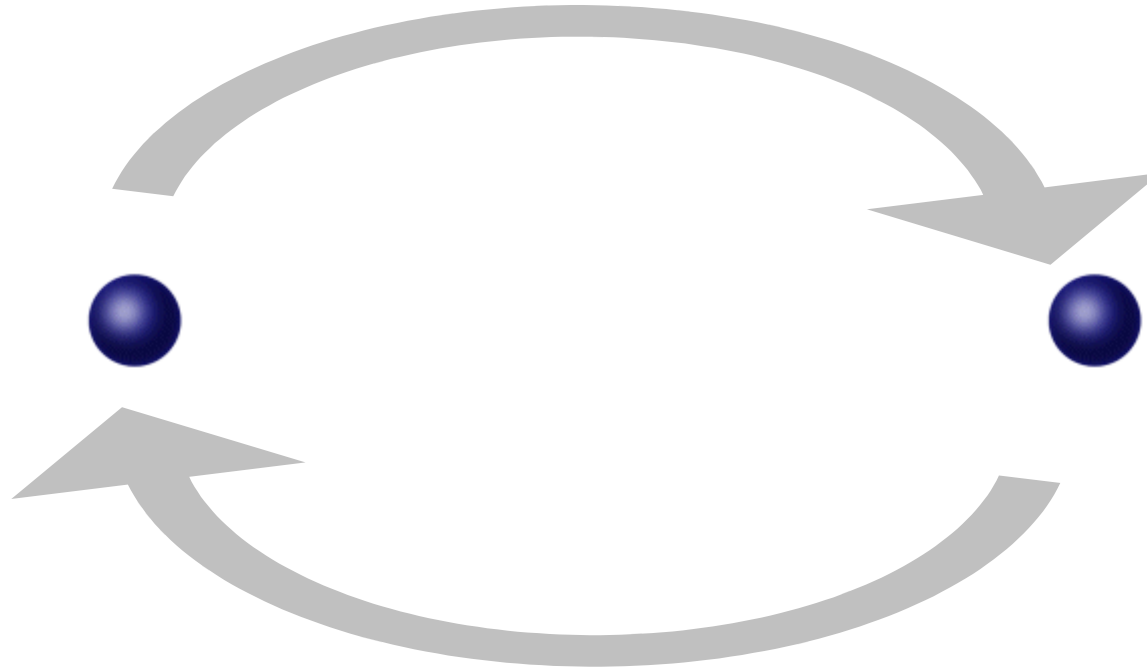
The topological order (TO) is defined by 3 features:

- the ground state degeneracy depends on the topology of space.
- there are anyonic quasiparticle excitations.
- there are edge excitations – a practical way to measure TO.



Anyons

Fermions and bosons



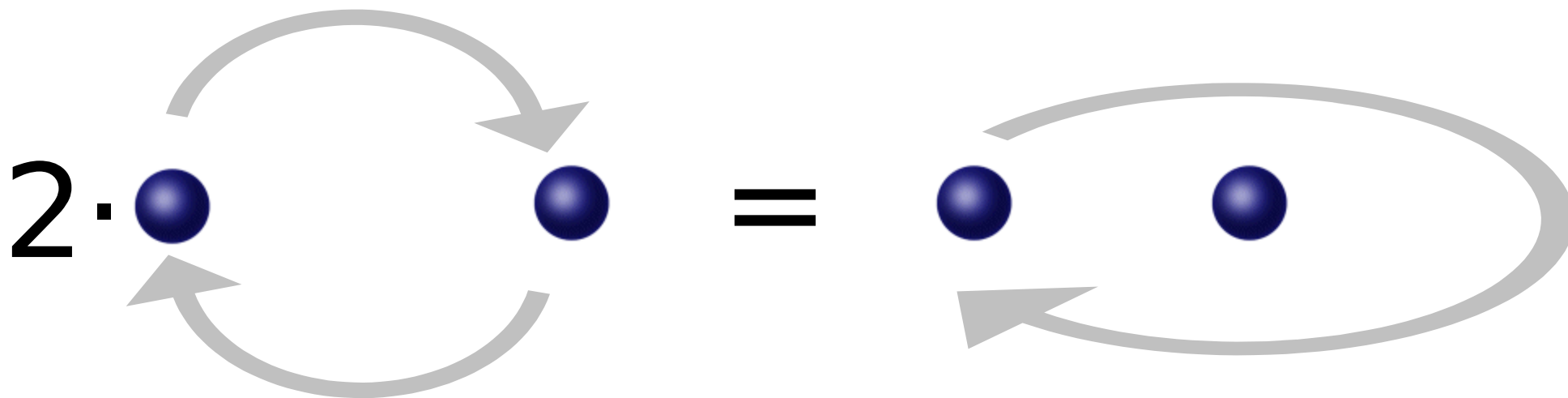
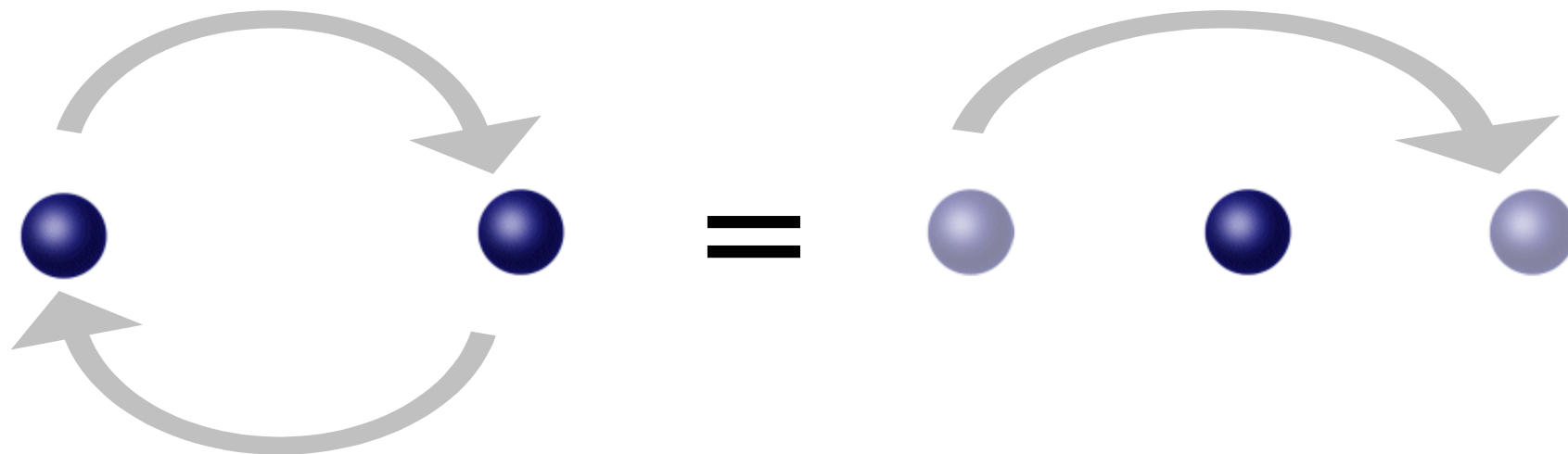
Bosons

$$|b_1 b_2\rangle = |b_2 b_1\rangle$$

Fermions

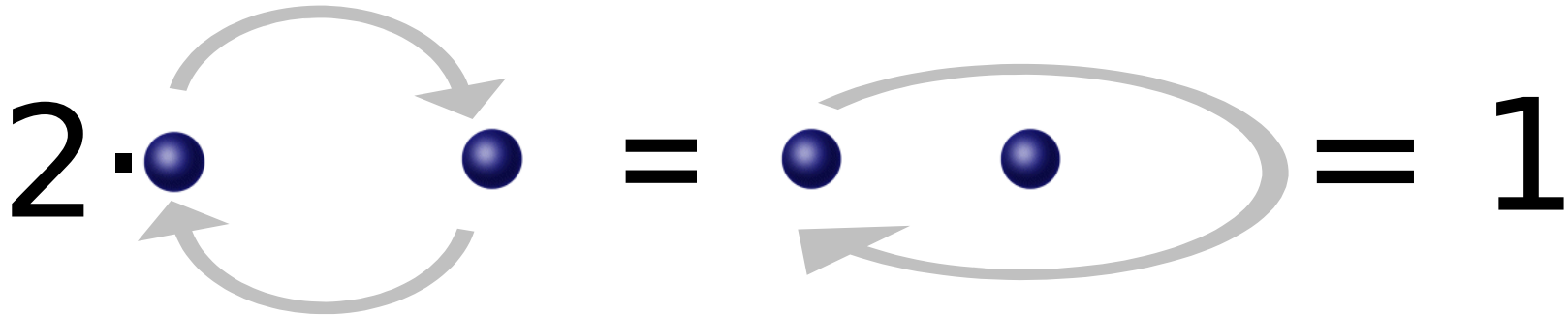
$$|f_1 f_2\rangle = -|f_2 f_1\rangle$$

Loops and interchanges

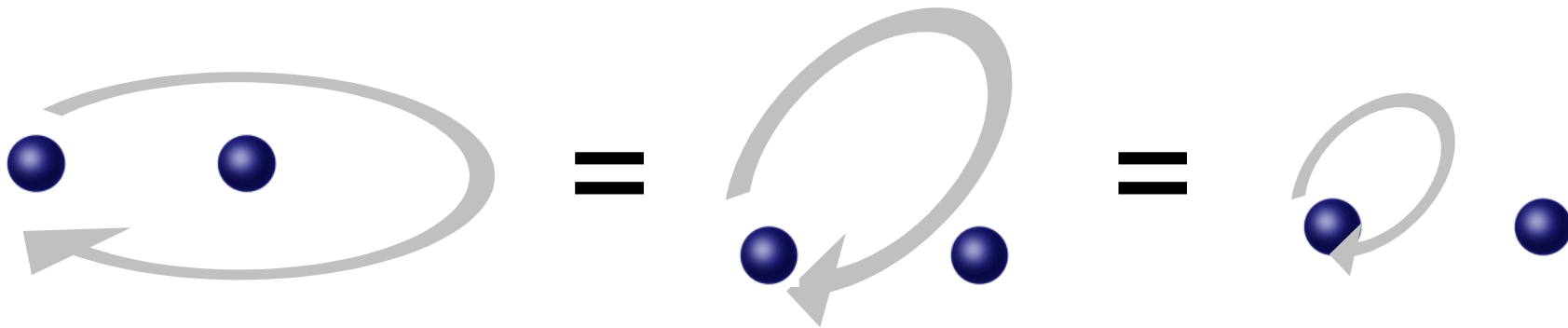


Loops and interchanges

In 3 dimensions:



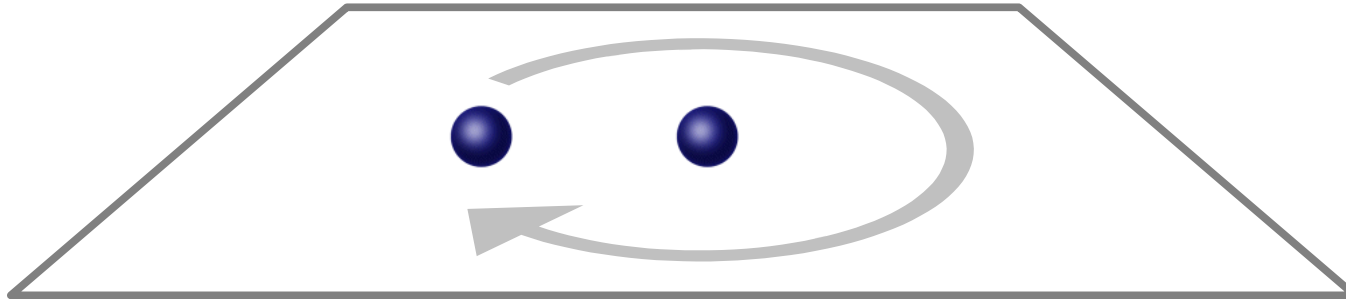
since we can **continuously** deform the loop around the other particle into the **identity**.



This is the reason why in **3 dimensions** we can only have **bosons** or **fermions**.

$$|p_1 p_2\rangle = e^{i\theta} |p_2 p_1\rangle \Rightarrow e^{i2\theta} |p_1 p_2\rangle = |p_1 p_2\rangle \Rightarrow e^{i\theta} = \pm 1$$

Nevertheless, in 2 dimensions we cannot contract the loop into the identity.



This makes possible the existence of particles that are neither bosons nor fermions. We will call them anyons.

$$|a_1 a_2\rangle = e^{i\theta} |a_2 a_1\rangle$$

Abelian and non-abelian anyons

Let us assume that our particles have an **internal degree of freedom** (colour, spin, etc.).

Then, the most **general operation** that the particles would suffer under an **interchange** would be,

$$|a_1^\alpha a_2^\beta\rangle = U_{\gamma\delta}^{\alpha\beta} |a_2^\gamma a_1^\delta\rangle$$

where $U_{\gamma\delta}^{\alpha\beta}$ is a **unitary** matrix.

Abelian anyons

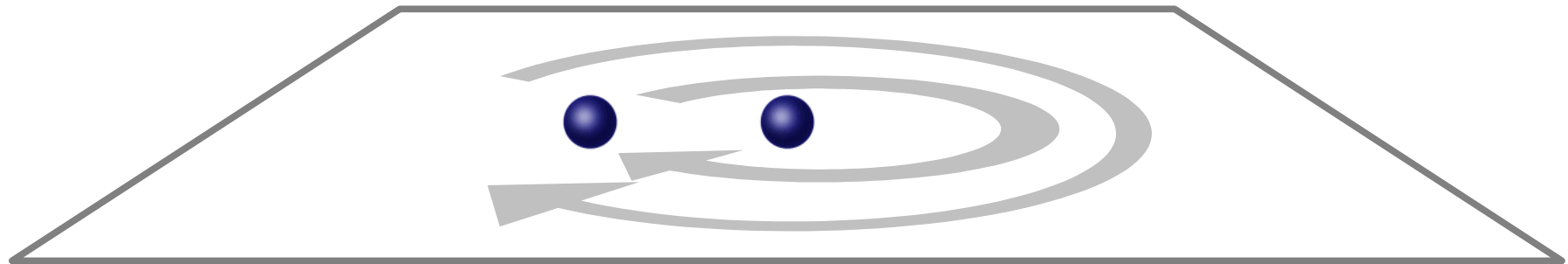
The generators of U belong to an abelian group

Non-abelian anyons

The generators of U do not belong to an abelian group

Applications of anyons

We have a **robust** method of performing **unitary operations** on quantum states.



$$|a_1^\alpha a_2^\beta\rangle = U_{\gamma\delta}^{\alpha\beta} |a_2^\gamma a_1^\delta\rangle$$

The final **unitary operation** performed does not depend on the particular **path** followed around the particle.

Can we take profit of this technology?

Topological quantum computation

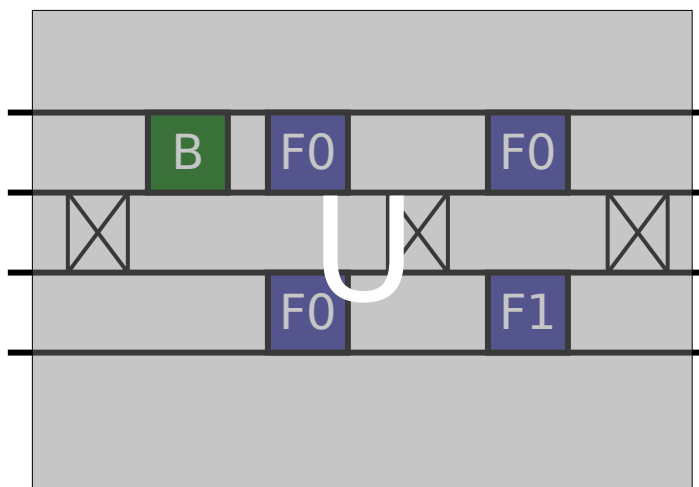
A quantum computer

A **quantum computer** is a computer that works with **qubits** instead of **bits**.

$$|\Psi\rangle = c_0|0\rangle + c_1|1\rangle$$

By means of **quantum gates** it can perform a **unitary operation** on its register of **qubits**.

$$|\Psi\rangle = \sum_{i_1 i_2 \dots i_n} c_{i_1 i_2 \dots i_n} |i_1 i_2 \dots i_n\rangle$$



- **Single** and **two qubit** gates.
- **Local** gates.
- **Polynomial** number of gates.

Quantum and classical computations

Quantum computer



>>

exponentially faster

Classical computer

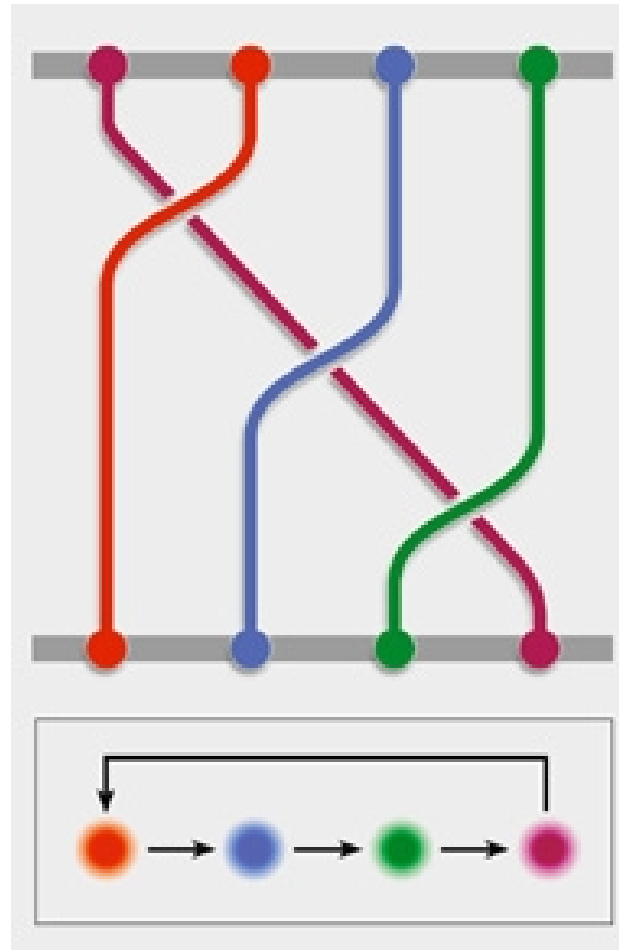


A quantum computer allows us to simulate quantum mechanics, to solve NP classical problems: factorization (Shor's algorithm), searches in data bases (Grover's algorithm), etc.

Problem: experimentally is very difficult to implement unitary transformations (precision problems, decoherence, etc.).

Idea: to use anyons

Anyons are **robust** against **local perturbations**, thus, they do not have the problems of the **standard quantum computation** paradigm.



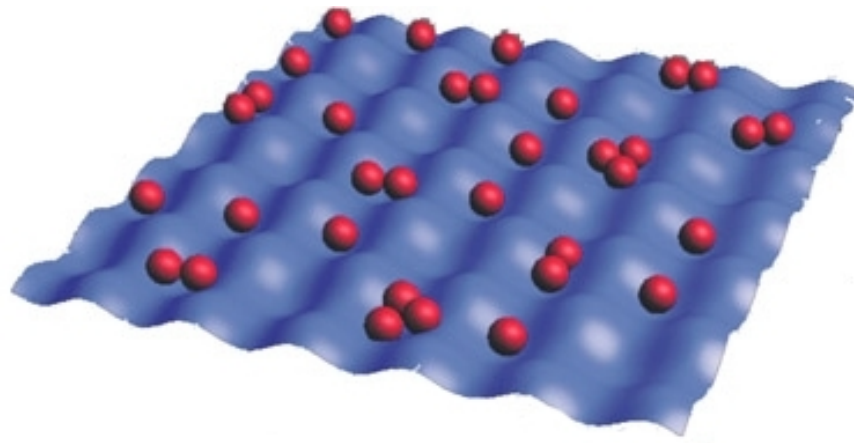
A **quantum computation** can be realized by braiding **anyons**.

Creating anyons

Anyons are collective excitations of the FQHE. Nevertheless, electrons are too small to be manipulated experimentally.

The solution to this problem is to realize another quantum system equivalent to the FQHE but accessible experimentally.

A suggestion is to use optical lattices, where the atoms would play the role of electrons and excitations could be created by means of an external laser.



Conclusions

- 1) The basic concepts related to the Landau symmetry breaking theory have been presented: phase transitions, order, symmetry, order parameter, gapless excitations, etc.
- 2) We have seen an explicit example of a new kind of order beyond the Landau's paradigm: the FQHE. This new kind of order is called topological order.
- 3) One of the signs of the topological order is the existence of the anyonic quasiparticles. They are collective excitations that are neither bosons nor fermions.
- 4) We have realized that anyons would be very useful in order to perform quantum computation. This is new paradigm of computation is called Topological Quantum computation.